

Geometric realisation of Connes spectral triples for algebras with central bases

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Geometric
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Preliminaries

NCRG motivation

NCRG formalism

NCRG formalism

Spectral
Triples

Axiomatic
formalism

Comparison with
Connes

Some results for
central bases

Some results for
central bases

Conclusions

1 Preliminaries

- NCRG motivation
- NCRG formalism
- NCRG formalism

2 Spectral Triples

- Axiomatic formalism
- Comparison with Connes
- Some results for central bases
- Some results for central bases

3 Conclusions

NCRG statement for quantum gravity

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- 1** We will not assume that the spacetime is continuum at the plank scale.
- 2** Instead propose that it is more effectively described by a noncommutative coordinate algebra.
- 3** And in the limit classical RG is recovered.

¹E.J. Beggs and S. Majid, Quantum Riemannian Geometry, Grundlehren der mathematischen Wissenschaften, Vol. 355, Springer (2020) 809pp 

Basic NCRG formalism

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We work with A a unital algebra, typically a \ast -algebra over \mathbb{C} , in the role of 'coordinate algebra'.

- 1 Differentials are formally introduced as a bimodule¹ of 1-forms equipped with a map $d : A \rightarrow A^1$ obeying the Leibniz rule $d(ab) = (da)b + adb$.
- 2 Assume this extends to an exterior algebra $(\wedge; d)$ with $d^2 = 0$ and d obeying the graded-Leibniz rule and $d(1) = 0$.
- 3 A quantum metric is $g \in \mathcal{K}(A^1)$ and a bimodule map inverse $(;) : \mathcal{K}(A^1) \rightarrow A$.
- 4 A bimodule connection on A^1 is $r : \mathcal{K}(A^1) \rightarrow A^1$ obeying:
 - 1 $r(a \cdot !) = a \cdot r ! + da \cdot !$,
 - 2 $r(! \cdot a) = (r !) \cdot a + (! \cdot da)$
 with a unique 'generalised braiding' bimodule map $\tau : \mathcal{K}(A^1) \otimes A^1 \rightarrow A^1 \otimes \mathcal{K}(A^1)$.

In recent years, the quantum Riemannian geometry was extended to a systematic theory including the QLC and further structure

- 1 'spinor' bimodule S equipped with a bimodule connection
- 2 A 'Clifford action': $\sigma^2 = 1$, $\sigma A \sigma = S$, $\sigma S \sigma = A$.
- 3 Leading to a quantum-geometric Dirac operator D .
- 4 And an inner product used to complete to a Hilbert Space.

Motivation for central bases

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spectral
triples for
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Triples

Axiomatic
formalism

Comparison with
Connes

Some results for
central bases

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We say a basis $\{e^j\}$ is central if it is a Grassmann algebra: $e^j e^j = 0$.

We will think in central bases A if A has trivial centre, S has a central basis $\{f^j\}$ and S has a central basis $\{g^j\}$.

The central bases assumption resumes into the set of 1-forms ω and its adjoint. The metric translates to the matrix of metric coefficients in the basis being hermitian. If the ip map then the ω -preserving condition translates to the Christoffel symbols in the basis being real.

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central bases

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NCRG motivation

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NCRG formalism

Spectral
Triples

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formalis9051 0.9196 0.9451 rg 0.9051 0.9196 0.9451 RG 0.39262 0.48549 0.64862 rg 0.39262 0.48549 0.64862 RG 0 -7.389 Td158 rg 0 0 0 0 8r2ompG /F5_0stric

Some results and examples

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realisation of
Connes
spectral
triples for
algebras with
central bases

Preliminaries
NCRG motivation
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Triples

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central bases

Some results for
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Theorem

Up to a phase in the 2D spinor bundle case, J can be obtained with $r > 0$ and $z \in \mathbb{C}$ as either:

$$(1): J = \begin{pmatrix} z & r \\ -\frac{z}{r} & -z \end{pmatrix}; \quad (2): J = \begin{pmatrix} 1 & \frac{-1}{|z|^2} z \\ z & -\frac{z}{z} \end{pmatrix}$$

or its transpose. The $|z| = 1$ case of (2) needs $z \neq 0$ and up to a

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formalism

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Noncommutative Torus

the only geometrically realised Dirac operator for the standard Euclidean metric and a WQLC are

$$D(e) = (\partial + s^j) \cdot e + d_j s^j \cdot e = \sum_i ((\partial + d_i) \cdot) e : (1)$$

. With Hilbert space, the state $\tau^m v^n = \sum_{m;0} \sum_{n;0} \mathbb{R}$



